

Features of discrete VLF emissions observed at Gulmarg, India during the magnetic storm of 6–7 March, 1986

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During the analysis of archived VLF data from Indian low latitude ground stations, some discrete VLF emissions recorded at the low latitude ground station Gulmarg (geomagnetic latitude 24°26'N; geomagnetic longitude 147°09'E, $L = 1.28$) during moderate magnetic storm activity ($\sum K_p = 32$, K_p index varies from 4 to 6 during the observation period) on 6/7 March, 1986 are presented in this paper. The dynamic spectra of these discrete VLF emissions were observed along with tweeks and its harmonics, which is interesting and complex to explain. In most of the events the harmonic frequency of tweeks correlates with the starting frequency of harmonics of discrete emissions. In order to explain the observed features of discrete VLF emissions, we propose cyclotron resonance interaction between whistler mode wave and energetic electrons of inner radiation belt as possible generation mechanism. An attempt is also made to determine parallel energy, anisotropy and wave growth relevant to the generation process of VLF emissions.

1. Introduction

Magnetospheric ELF/VLF emissions is a class of natural radio phenomena, which is often observed in close association with whistlers both at the Earth's surface and onboard satellites. They are grouped as: (a) unstructured continuous emissions in both time and frequency which tend to maintain a steady state like hiss, resonance band and noise bands near the ion gyrofrequencies, and (b) structured discrete emissions with a repetitive and even periodic character which tend to be transient like chorus, periodic emissions and other transient discrete emissions such as hooks, risers, fallers, pseudo whistlers (Helliwell 1967; Hattori *et al* 1991; Sazhin and Hayakawa 1992; Bell *et al* 2000). The generation of discrete VLF emissions is one of the most puzzling problems of VLF waves in the Earth's magnetosphere. Although it is generally believed that their generation is

connected with the cyclotron resonance of whistler-mode waves and radiation belt electrons (Helliwell 1967), the mechanism responsible for the origin of discrete VLF emissions and the formation of spectrum of separate elements are still a subject of active experimental and theoretical research (Trakhtengerts 1999; Santolik and Gurnett 2003; Santolik *et al* 2004; Singh and Singh 2004 with references therein). At mid- and high-latitudes the observed correlation between discrete VLF emissions and energetic electrons suggests the key role of the latter in the generation of the former. Nunn and Sazhin (1991) considered interaction between VLF hiss and energetic electrons and explained the similarities between fine structure of chorus and hiss emissions. To explain the spectrogram of hiss and discrete VLF emissions, Dowden (1971) suggested that whistlers echoing in the discrete emission path are initiated by sferics and appear to be the source of discrete VLF emissions and

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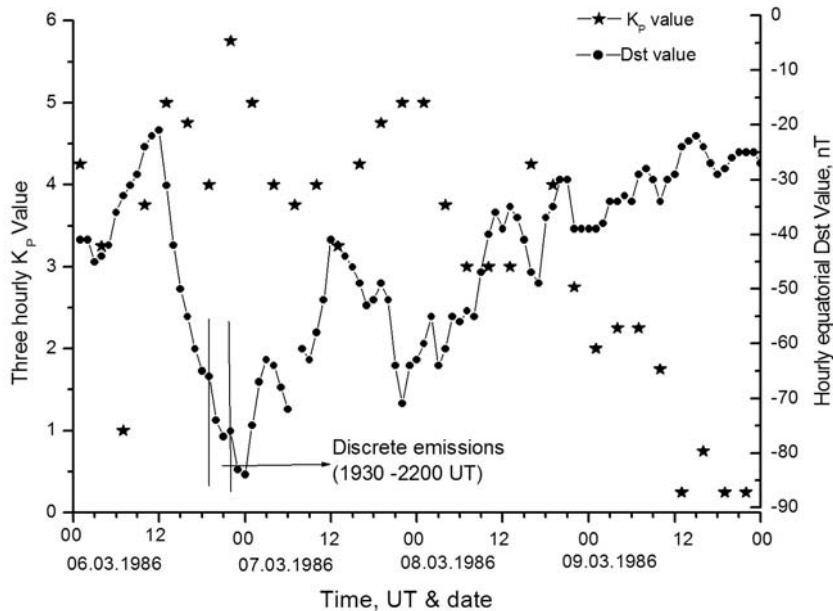


Figure 1. Equatorial hourly Dst variations and three hourly K_p index during the main phase magnetic storm 6–9 March, 1986. The emission period of the event is also marked.

hiss bursts involving lightning triggered from the Earth's magnetosphere.

Since the propagation characteristics of ELF/VLF emissions in the low latitude ionosphere are not properly known, the mechanism of their generation source and propagation are far from being well understood (Singh *et al* 2000; Singh and Singh 2004). Therefore, a better understanding of the generation mechanism for ELF/VLF emissions observed at low latitudes would be useful for analyzing the properties of high energy trapped electrons. During the course of our analysis of old ELF/VLF data recorded during January 1986 to July 1986 at low latitude ground station Gulmarg (geomag. lat. $24^{\circ}26'N$; geomag. long. $147^{\circ}09'E$, $L = 1.28$), we found a new type of discrete VLF emission which we report here with a discussion of their probable generation mechanism.

In the present paper, first we present spectral analysis of the ELF/VLF discrete emission events observed at Gulmarg during the routine recording of whistlers. The dynamic spectra of the discrete VLF emissions were observed along with tweaks and its harmonics. In most of the events, the lowest frequency of harmonics of discrete emissions correlates with the second and third harmonic frequency of tweaks. Further, generation mechanism of these discrete VLF emissions has been proposed and an attempt is made to determine parallel energy, anisotropy and wave growth relevant to the generation process of discrete emissions. Finally results are discussed with other published results.

2. Experimental observations and analysis

At the low-latitude ground station Gulmarg, the wideband ELF/VLF waves were received by a T-type antenna, suitably amplified by pre- and main-amplifiers and recorded using a tape recorder. The recorded data were analysed by digital Sonograph machine and Advance VLF Data Analysis System (AVDAS). At low latitudes, nights with observable ELF/VLF emissions are rather rare, and the activity is closely related to strong magnetic activity (Singh *et al* 2005). Several interesting ELF/VLF events were recorded during a magnetic storm often large enough to allow for a statistical analysis. The observations of discrete VLF emissions at low-latitude ground station Gulmarg are unusual in the sense that most of the reported discrete emissions are from satellites (Cornilleau-Wehrin *et al* 1978; Hattori *et al* 1991; Santolik and Gurnett 2003; Santolik *et al* 2004). In this paper, the discrete VLF emissions recorded during the night of 6–8 March, 1986 are analysed. We have recorded a new type of discrete VLF emission, which is associated with the tweaks and its harmonics. These discrete VLF emissions were recorded during the magnetic storm period of 6–9 March, 1986, with minimum Dst index of -84 nT on 6 March and maximum $\sum K_p = 34$ on 7 March. The Dst -index variations and K_p -index variations are shown in figure 1, in which the observation period of the event is also marked. It is to be noted that these discrete events were observed during the main phase of the magnetic storm.

07 March, 1986 GULMARG
Geomag. Lat. $24^{\circ}26' N$, Long. $147^{\circ}09' E$

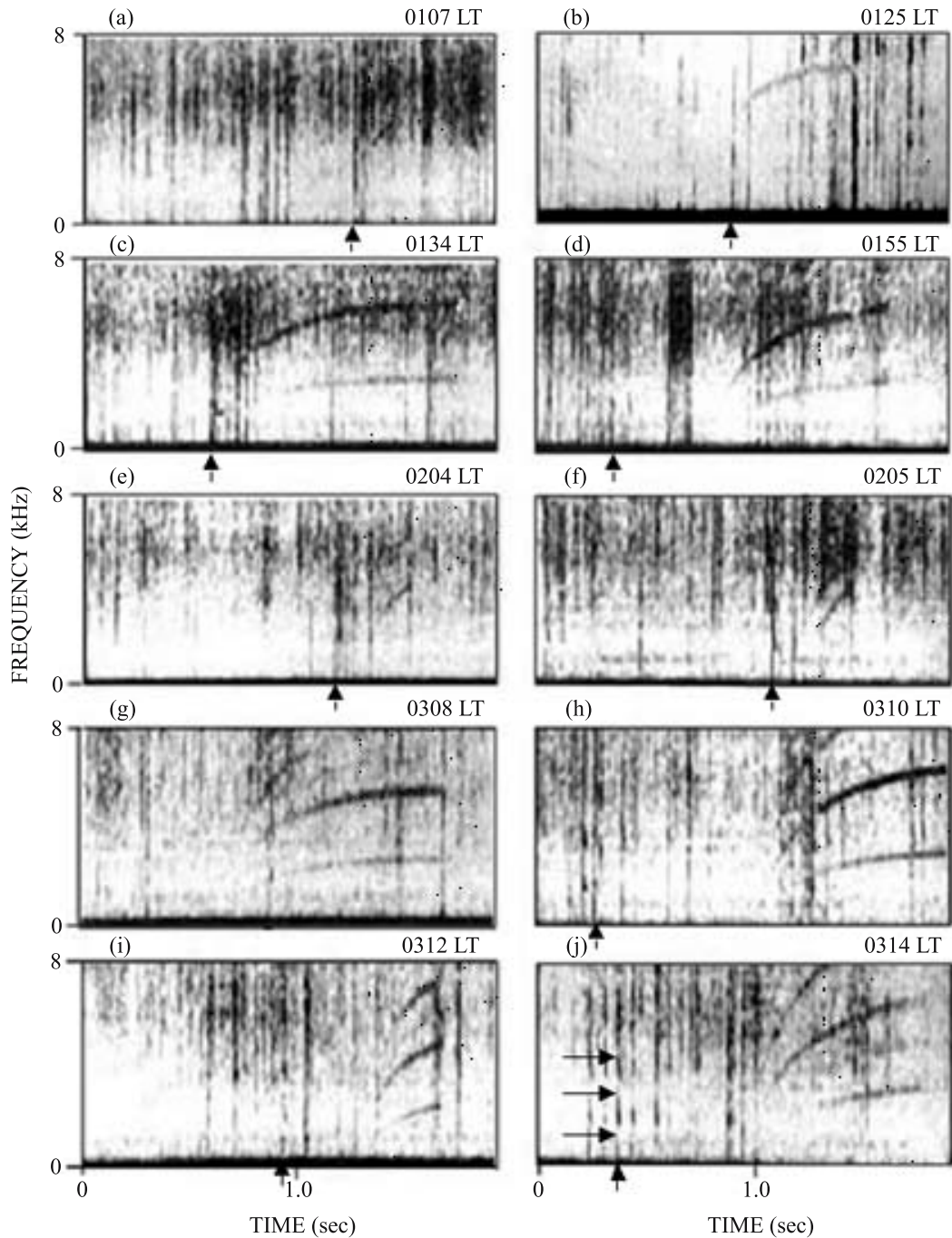


Figure 2. Typical examples of lightning generated sferics along with discrete VLF emissions observed at Gulmarg on 7 March, 1986 (local time). Arrows mark the tweeks. The first, second and third harmonics of the tweeks are also marked in (j) by arrows on frequency scale.

2.1 Spectral analysis of discrete VLF emissions

Typical frequency-time spectrograms of sferics and discrete VLF emissions recorded at Gulmarg during the night hours of 6–7 March, 1986 are shown in figure 2. Discrete emission events started at

1930 UT (0100 LT) and continued up to 2200 UT (0330 LT). Figure 2 shows an association between VLF sferics (tweeks) and discrete VLF emissions. Here we note that for most of the events the fundamental frequency of tweeks corresponds to the initiating frequency of fundamental discrete emissions. The second harmonic frequency of tweeks

corresponds to the starting frequency of the second harmonics of discrete emissions and so on. Tweaks are ELF/VLF electromagnetic waves which originated in the distant lightning and propagated in the Earth-ionosphere waveguide. They are known to exhibit remarkable dispersion near the cut-off frequencies of the first-, second-, third-order mode waves, and these dispersion effects are found to fit well with the Earth-ionosphere waveguide propagation theory (Singh and Singh 1996). In the course of time the upper boundary frequency of discrete VLF emissions was fluctuating and the number of events increased with the passage of time. First harmonic tweaks (sferics) accompanied with the single event of discrete VLF emissions observed in different frequency ranges between 2.6 and 6.5 kHz is shown in figure 2(a, b and e), whereas figure 2(c, d and f) contain second harmonic tweaks and correspondingly two discrete VLF emission events lying in different frequency ranges between 2 and 6.5 kHz. Figure 2(g, h, i and j) clearly show third harmonic tweaks accompanied by three events of discrete VLF emissions of short and long duration in different frequency ranges between 1 and 8 kHz. Hence, a fairly consistent pattern emerges from the observed features of the tweaks and discrete VLF emissions recorded at Gulmarg and a correlation is found between them.

From figure 2, an interesting point emerges concerning the correspondence between the cut-off frequencies of harmonics of tweaks and starting frequency of corresponding discrete emissions. For example, figure 2(j) shows that the first harmonic frequency is 0.7 kHz and the initiating of discrete emission is 0.7 kHz. Similarly the cut-off frequency of the second harmonic is 1.9 kHz and corresponding starting frequency of the second discrete VLF emission is 1.9 kHz. Corresponding frequency for the third harmonic and the third discrete emission are 3.4 kHz and 3.4 kHz respectively. The observed discrete VLF emission elements have the following mean parameters: $f_{\min} = 1$ kHz, $f_{UB} = f_{\max} = 5.04$ kHz and $df/dt = 6.622$ Hz sec⁻¹. The standard deviation of f_{UB} was found to be about 0.65. The discrete VLF emissions shown in figure 2 occurred in the wide frequency range between 1 and 8 kHz with the rate of change of frequency with time (df/dt) between 0.9 and 10.5 kHz sec⁻¹. The relationship between the cut-off of the tweak and the cut-off of the VLF emissions indicates that the emissions may have actually propagated in the Earth-ionosphere wave-guide and this correlation in frequency could be due to wave-guide cut-off.

3. Generation mechanism

The dynamic spectrum of tweaks (sferics) and its higher harmonics can be understood by

considering part of VLF energy after lightning discharge to propagate through the Earth-ionosphere waveguide (Singh and Singh 1996; Shvets and Hayakawa 1998). Discrete VLF emissions can be explained by transverse resonance interaction between whistler mode waves and counter streaming energetic electrons (Helliwell 1967; Nunn and Sazhin 1991). Analyzing Ariel 3 and 4 satellite data Hayakawa (1989) has also suggested that low latitude VLF emissions may have originated during lightning discharges.

The observed features of tweaks shown in figure 2 and cut off frequency of tweaks confirm that the scenario in which lightning source is closer to the observation point in the northern hemisphere may give rise to higher harmonics. Usually lightning energy is coupled to the Earth-ionosphere wave-guide and tweaks along with the higher harmonics are generated. Only in rare cases, lightning generated electromagnetic energy propagates both along the geomagnetic field lines and in the Earth-ionosphere waveguide, for which tweaks are found to be correlated with VLF whistler-mode waves. The second and the third harmonic of discrete VLF emissions may have been generated through the wave-particle interactions in the magnetosphere. The two events having their origin in the lightning discharge and wave-particle interaction respectively are very well correlated.

As the wave is received at the low latitude Earth's station, Gulmarg, we may consider that the generation region could be near the equatorial plane of the field line corresponding to Gulmarg in the inner zone radiation belt ($L \sim 1.2$). In this region there are a large number of energetic electrons, which can effectively participate in the generation mechanism through Cyclotron resonance interaction (Rycroft 1972; Imhoff *et al* 1973). In order to explain the observed features of discrete VLF emissions we propose cyclotron resonance interaction between whistler mode wave and energetic electrons of inner radiation belt. We note that for effective transfer of energy from the interacting particle to the wave, resonance condition has to be satisfied which is written as

$$1 - \beta_{\parallel} = \frac{\omega_H}{\gamma\omega}, \quad (1)$$

where $\beta_{\parallel} = v_{\parallel}/c$, v_{\parallel} is the electron velocity along the magnetic field and $\gamma = (1 - \beta_{\parallel}^2)^{-1/2}$. The refractive index for parallel propagating waves under condition $\omega^2 \ll \omega_H^- \omega_H^+$ is much greater than 1 and is written as (Stix 1962).

$$n^2 = \frac{\omega_{pe}^2}{\omega_{He}(\omega + \omega_{Hi})}, \quad (2)$$

where $\omega_{pe}\omega_{He}$ are the electron plasma frequency and electron gyrofrequency respectively. ω_{Hi} is the proton gyrofrequency. The resonant energies E_{\parallel} for various frequencies of emissions can be written as (Tsurutani *et al* 1975)

$$E_{\parallel} = \left(\frac{\omega_{He}}{\omega_{pe}}\right)^2 \left(\frac{\omega_{He}}{\omega}\right) \left(1 + \frac{\omega_{Hi}}{\omega}\right) m_0 c^2, \quad (3)$$

where m_0 is the rest mass of electron, c is the velocity of light in vacuum.

The interacting waves can be amplified if the high-energy tail of the velocity distribution function has some finite anisotropy $A = (T_{\perp}/T_{\parallel}) - 1$, where T_{\perp} and T_{\parallel} are the temperatures of the electrons perpendicular to and parallel to the geomagnetic fields respectively. In the regime of linear theory, the growth rate of the wave for $\omega \ll \omega_{He}$ is given by (Kennel and Pestschek 1966)

$$\Gamma = \pi\omega_{He} \left(1 - \frac{\omega}{\omega_{He}}\right)^2 \frac{J(> E_R)}{2V_R N_T} \times \left[A(V_R) - \frac{1}{(\omega_{He}/\omega) - 1} \right], \quad (4)$$

where $J(> E_R)$ is the omnidirectional flux of electrons having energy greater than the resonance energy. V_R is the resonance velocity and N_T is the total number density of electrons.

4. Results and discussions

The observation of ELF/VLF emissions at low-latitude station Gulmarg shows that, although discrete VLF emissions are rare under quiet conditions the occurrence rate of discrete VLF emissions is high during magnetic storms. We have recorded a new type of discrete emission during the magnetic storm period of 6–9 March, 1986. During substorms the inflow of the particles with energy of about 1 keV from the plasma sheet into the inner magnetosphere takes place. Plasma is trapped by the geomagnetic field and accelerated. The oscillating particles drift to dawn (electrons) and to dusk (protons). The large-scale dawn-dusk electric field, which is increased during disturbed conditions cause plasma to drift across L shells. As electrons drift to dawn the conditions become more and more favourable for the cyclotron instability to arise (Bespalov and Trakhtengerts 1986) and the background plasma density increases from midnight to morning-noon hours owing to filling of the equatorial regions of the magnetosphere with cold ionospheric plasma. At a certain time the energetic electrons reach a certain region where the conditions for the cyclotron instability turn out to be

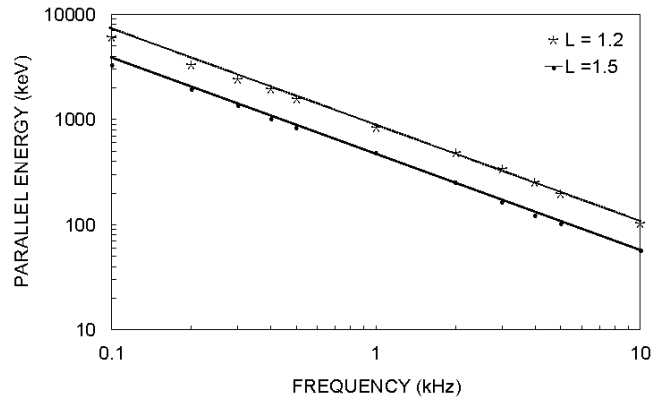


Figure 3. The variation of parallel energy of resonating electrons with wave frequency at L = 1.2 and 1.5.

optimal. In this region the dissipation of accumulated energy takes place by cyclotron instability. As a result of this dissipation the generation of ELF/VLF emissions and precipitation of resonant electrons into the loss-cone proceeds (Smirnova 1984).

In order to test cyclotron resonance interaction as a possible generation mechanism for discrete VLF emissions, we have computed the resonance energy of the high-energy interacting electrons and growth rate of whistler waves at L = 1.2 in the equatorial plane. The variations of parallel energy of resonating electrons with wave frequency for different L-values are shown in figure 3. It is found that the energy of resonating electrons decreases as L-value increases for a given frequency band. Burton and Holzer (1974) have shown that the discrete emissions are generated by cyclotron resonance with electrons in the energy range of 5–150 keV with pitch angle distribution peaked at 90° to \mathbf{B} and anisotropy greater than a critical value. Further, it has been shown that the resonant energies for various frequencies of the emissions at L = 1.2 are in the MeV range (Lalmani *et al* 1972). Thus, the computed values of resonant energy for various frequencies are in good agreement with the reported results (Lalmani *et al* 1972; Burton and Holzer 1974).

Equation (4) shows that the wave growth linearly increases with anisotropy A in the velocity distribution function. Further Γ is positive only when $A > \omega/(\omega_{He} - \omega)$. Thus there is a minimum anisotropy which is required for the wave to grow and hence for the generation of ELF/VLF emissions. The variations of this minimum anisotropy as a function of wave frequency for different L-values are shown in figure 4. The anisotropy A increases with frequency as well as L-value. It is also seen that the gyroresonance leading to extremely low frequency signals requires only small

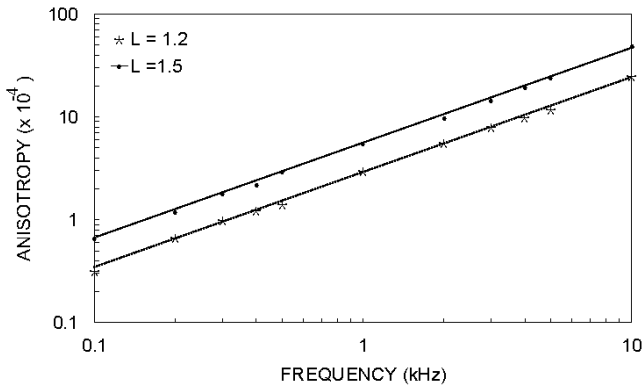


Figure 4. The variation of the minimum anisotropy as a function of wave frequency at $L = 1.2$ and 1.5 .

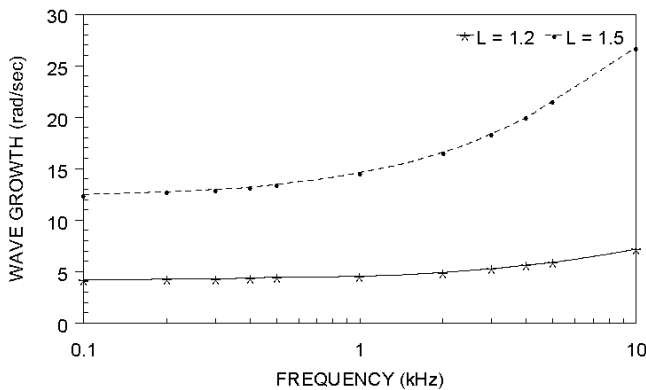


Figure 5. The variation of wave growth rate with excited frequency at $L = 1.2$ and 1.5 .

values of A , however on a particular field line, this requires resonance with electrons of much higher energy.

For computation of wave growth, the density of energetic electrons at different L -values are derived from the measurement of Katz (1966) who has reported the variation of electron flux as a function of energy for different L -values in the inner zone radiation belt. The variation of wave growth rate with excited frequency for $L = 1.2$ and $L = 1.5$ is shown in figure 5. It is observed that the growth rate is larger and it increases with frequency. This large amplification causes significant enhancement in wave amplitude, which in turn triggers discrete emissions to be observed at low latitude ground station Gulmarg.

The rising frequency spectrum of discrete VLF emissions observed at Gulmarg can be explained by considering the interaction region to start from the equator and extend to some finite length in the southern hemisphere along the geomagnetic field line. After moving away from the equator, the local electron gyrofrequency becomes too large for the resonance condition to be still met and the

generation of rising tone ceases. Thus, the maximum frequency of discrete emissions is controlled by the extension of interaction region along the geomagnetic field line.

In the above, we considered the generation of discrete emission through the process of wave-particle interaction in the magnetosphere. However, there is no direct evidence that these emissions belong to the magnetospheric origin. These emissions could be generated during lightning discharge along with tweeks. In this case it is not clear how to explain dispersion produced in the VLF signal called discrete emissions. The other possibility is that discrete emission is generated during instability in the ionosphere. Again this aspect has not been explored in detail.

5. Conclusion

A detailed spectral analysis of a new type of discrete VLF emissions observed at the low-latitude ground station Gulmarg during the strong magnetic activity period have been carried out. The possible generation mechanism for temporal and spectral features of these discrete VLF emissions is presented. It is suggested that these discrete emissions could be generated during resonant cyclotron interaction in the equatorial zone of the inner magnetosphere. A further experimental and theoretical study of discrete VLF emissions at low latitude would definitely contribute to a more detailed understanding of these phenomena.

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